

# Calculation of Stored Electromagnetic Powers and Q Factors of Very Small Normal-Mode Helical Antennas

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**Abstract**— Theoretical explanations of radiated, dissipated and stored energies of small antennas were shown in a textbook. However, practical values of these energies were not discussed previously. In this paper, these energies are obtained numerically through electromagnetic simulations. As for a study object, a normal-mode helical antenna is utilized, because stored electric and magnetic energies are simultaneously observed. First of all, it is shown that these two energies have the same value at the self-resonant conditions. Next, electric and magnetic stored powers are calculated from these energies and are compared with the input power. Moreover, antenna *Q* factors are also obtained. Adequateness of calculated results is ensured thorough comparing with other calculated results.

## I. INTRODUCTION

For designing high efficiencies at small antennas, to understand effects of stored electric and magnetic energies is important. Stored electric and magnetic energies and dissipated energies of antennas were detailed discussed in the Harrington's book [1]. However, actual values of individual energies were not shown before. Now, electromagnetic simulations become reliable and are expected to be useful for clarifying above mentioned energies.

In this paper, normal-mode helical antennas (NMHA) are selected for study objects, because stored electric and magnetic energies can be observed simultaneously [2]. First of all, antenna structures used in this study are explained with respect to self-resonant structures. And calculation process of stored electric and magnetic energies is shown. Next, stored energies are converted to stored powers and compared with the input power. Furthermore, antenna *Q* factors are obtained. And, *Q* factors obtained from three methods are compared and adequateness is ensured.

## II. SELF-RESONANT STRUCTURES

Antenna structure of a normal-mode helical antenna is shown in Fig. 1.  $H_A$ ,  $D_A$  and  $d$  express antenna height, diameter and wire diameter, respectively.  $a$  expresses the radius of a sphere enclosing the antenna. The feature of this antenna is to have two radiation sources such as very small electric current and magnetic current sources simultaneously [3]. And the electric source produces electric fields ( $E$ ) and the magnetic source produces magnetic fields ( $H$ ), respectively. The input impedance ( $Z_{in}$ ) is given by the next expression.

$$Z_{in} = R_{rad} + R_{metal} + j(-X_C + X_L) \quad (1)$$

Here,  $R_{rad}$  and  $R_{metal}$  indicate radiation and metallic loss resistances, respectively. And,  $R_{in}$  is used for  $R_{rad} + R_{metal}$ .

$X_C$  and  $X_L$  indicate input capacitance and inductance, respectively. At  $X_C = X_L$ , the reactance part becomes zero. This situation is called the self-resonance. At this condition, stored electric and magnetic powers are supposed to be cancelled out. And all the input power is effectively utilized for the radiation.

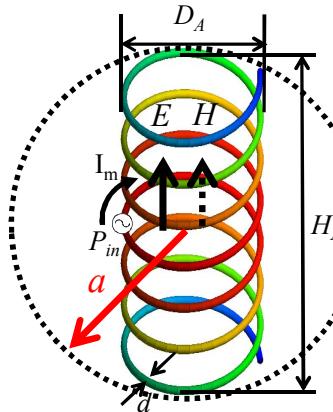


Fig. 1 Normal-mode helical antenna (NMHA)

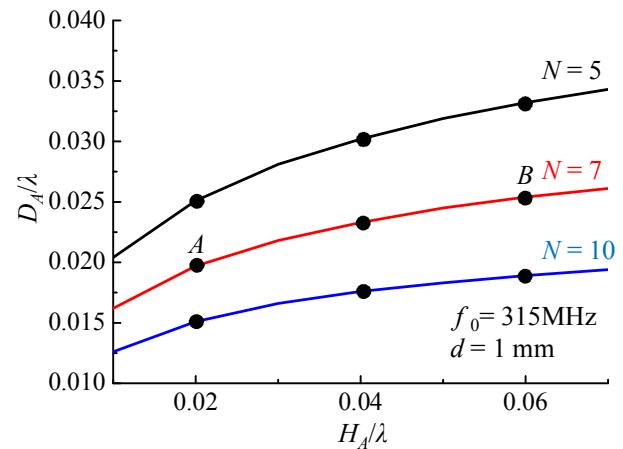


Fig. 2 Self-resonant structures

The self-resonant structures are determined through the next structural equation [2].

$$600\pi \frac{19.7N(\frac{D_A}{\lambda})^2}{9\frac{D_A}{\lambda} + 20\frac{H_A}{\lambda}} = \frac{279\frac{H_A}{\lambda}}{N\pi(0.92\frac{H_A}{\lambda} + \frac{D_A}{\lambda})^2} \quad (2)$$

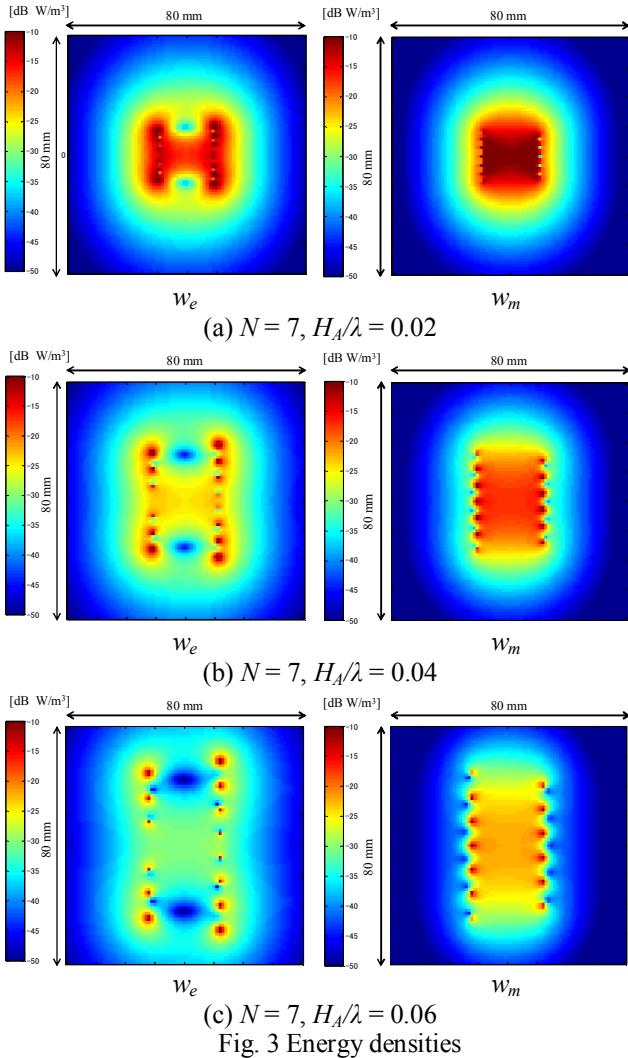
The self-resonant structures are shown in Fig. 2.  $D_A$  and  $H_A$  exist on a curve having a parameter  $N$  (number of turns). From the inclination of curves,  $D_A$  changes are rather small for the  $H_A$  changes. At the black circle points, simulation data are obtained.

### III. STORED ELECTRIC AND MAGNETIC ENERGIES

Calculation process of stored electro and magnetic energies and powers are explained in this section. And some discussions on comparisons of input and store powers are made. Moreover, stored power dependences on antenna lengths are considered.

#### A. Energy density distributions

Electric ( $w_e$ ) and magnetic ( $w_m$ ) energy densities are obtained based on the calculated results of electric ( $E$ ) and magnetic ( $H$ ) fields around the NMHA with the following equations.



$$w_e = \frac{\epsilon}{2} E^2 \quad (3)$$

$$w_m = \frac{\mu}{2} H^2 \quad (4)$$

Calculated results of  $w_e$  and  $w_m$  are shown in Fig. 3. Here, antenna structures of  $N = 7$  in Fig. 2 are used. And the input powers to the antennas are set 1 W. In the calculations, the input impedance miss matches are ignored. Thus, all antennas are receiving 1 W input power. At structures of small  $H_A$  values,  $w_e$  and  $w_m$  become very strong. The largest values of  $w_m$  at  $H_A/\lambda = 0.02, 0.04$  and  $0.06$  are  $-10$  dBW/m<sup>3</sup>,  $-17$  dBW/m<sup>3</sup> and  $-24$  dBW/m<sup>3</sup>, respectively.

The current amplitudes at the feed point ( $I_m$ ) of structures in Fig. 2 are shown in Fig. 4.  $I_m$  values increase rapidly in accordance with decreases of  $H_A$ . The input resistances ( $R_{in}$ ) and radiation resistances ( $R_f$ ) are shown in Fig. 5.  $R_{in}$  values decrease in accordance with decreases of  $H_A$ .

From the results of  $I_m$  and  $R_{in}$ , the antenna input power can be calculated with the next equation.

$$R_{in} \times (I_m)^2 / 2 = P_{in} [\text{W}] \quad (5)$$

All the results of Fig. 4 and Fig. 5 satisfy the results of  $P_{in} = 1$  W.

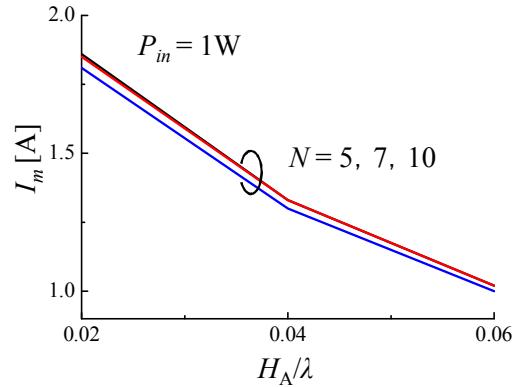


Fig. 4 Feed current

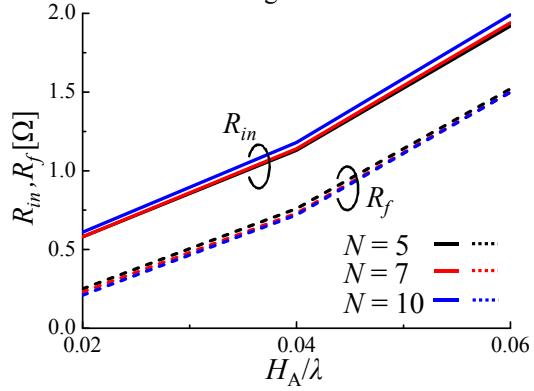


Fig. 5 Input impedance and radiation resistance

### B. Stored electric and magnetic powers

By employing the results of Fig. 3, stored electric and magnetic energies are calculated through the following equations.

$$W_e = \int_V w_e dv = \frac{\epsilon}{2} \int_V E^2 dv \quad (6)$$

$$W_m = \int_V w_m dv = \frac{\mu}{2} \int_V H^2 dv \quad (7)$$

As for the integration areas of  $V$ , the cubic of 80 mm lengths are used. The electric powers are estimated by the next expression [1].

$$P_{in} = P_f + P_d + j2\omega(W_m - W_e) \quad (8)$$

Here,  $P_{in}$  indicates the input power.  $P_f$  and  $P_d$  indicate radiated and dissipated powers, respectively.

As for stored powers, electric and magnetic powers are calculated by the following expressions.

$$P_e = 2\omega W_e \quad (9)$$

$$P_m = 2\omega W_m \quad (10)$$

Calculated results of  $P_e$  and  $P_m$  are shown in Fig. 6 [4]. In the case of integral calculation shown in Eqs.(6),(7), thin cylindrical area containing the antenna wire is omitted because irregular values appeared at  $E$  fields. Therefore,  $P_e$  values become somewhat smaller than  $P_m$ . By taking into account this calculation approximation, Good agreements of  $P_e$  and  $P_m$  are observed. From these agreements, the stored  $P_e$  and  $P_m$  cancel each other in Eq. (8). It is ensured that stored electric and magnetic powers are cancelled out at self-resonant structures. Moreover, values of  $P_e$  and  $P_m$  become between 500 W and 2000 W. These values are very large compare with the input power of 1 W. It is understand that the balanced situation is very critical. In very small deviations of structures, the balanced conditions are broken. These critical conditions imagine very narrow bandwidths.

Other features of  $P_e$  and  $P_m$  are dependences on  $N$  and  $H_A$ . These dependencies are discussed in the next section.

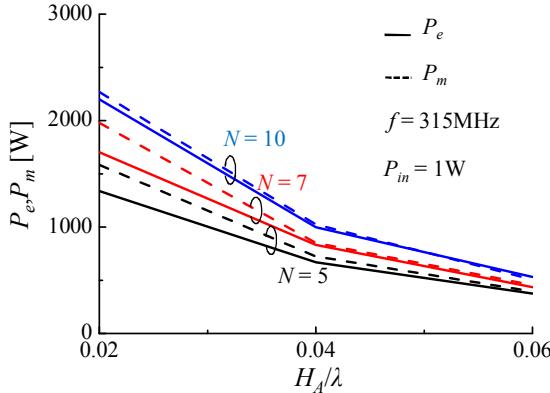


Fig. 6 Stored electric and magnetic powers

### C. $P_e$ and $P_m$ dependences on $N$ and $H_A$

In considering  $P_e$  and  $P_m$  values, expressions of Eq. (6) and Eq. (7) give the basic concepts.

As for  $P_m$ , magnetic field ( $H$ ) changes are important. Magnetic fields shown in Fig. 7 are calculated by the next expression.

$$H = \int_l \frac{I(l)}{4\pi r^2} dl \quad (11)$$

Here,  $I(l)$  indicates current on the antenna wire. Therefore,  $H$  becomes proportional to the input current ( $I_m$ ). Namely,  $P_m$  becomes proportional to  $I_m^2$ . Considering the tendencies of Fig. 4,  $P_m$  increases are recognized. Dependencies on  $N$  are anticipated from the relation of  $H$  to  $N$  as shown in Fig. 7.  $H$  is increased in accordance with  $N$  increases.

As for  $P_e$ , electric field ( $E$ ) changes are important. As shown in Fig. 8,  $E$  is determined by charges ( $Q$ ) at both ends. Then, the next proportion relation is given.

$$E \propto Q \quad (12)$$

At the same time,  $Q$  is related to the input current ( $I_m$ ) as shown in the next expression.

$$Im = \frac{dQ}{dt} = j\omega Q \quad (13)$$

As a result,  $E$  becomes proportional to  $I_m$ . Therefore,  $P_e$  becomes proportional to  $I_m^2$ . Considering the tendencies of Fig. 4,  $P_e$  increases are recognized.

Dependencies on  $N$  are recognized by comparing results of Fig. 8 (a) and (b). At larger  $N$  number, the cross sectional area containing  $E$  fields becomes small. Even if stored charge  $Q$  is unchanged,  $E$  fields inside  $N = 10$  become larger than the  $N = 5$  case. This tendencies explain the  $N$  dependence of  $P_e$ .

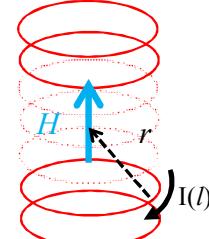


Fig. 7 Magnetic field expressions

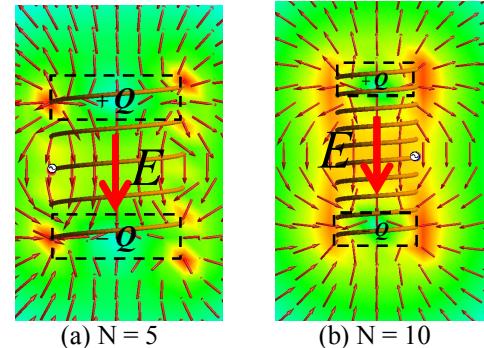


Fig. 8 Electric field expressions at  $HA/\lambda = 0.04$

#### IV. ANTENNA Q FACTORS

In the previous section, stored powers ( $P_{st}$ ) of electric and magnetic fields are numerically clarified. Then, the antenna  $Q$  factor can be calculated by the next expression.

$$Q = \frac{P_{st}}{P_f} \text{ or } \frac{P_{st}}{P_f + P_d} \quad (14)$$

Here,  $P_{st}$  represents  $P_e$  or  $P_m$ . And,  $P_f$  indicates the radiated power in Eq. (8).  $P_f$  and  $P_d$  are calculated by employing the results of Fig. 4 and Fig. 5.  $P_f$  is calculated by  $R_f \times I_m^2/2$ .  $P_d$  is calculated by  $(R_{in} - R_f) \times I_m^2/2$ .  $P_{st}/P_f$  is sometimes referred as an antenna  $Q$ . However, an actual antenna  $Q$  is given by  $P_{st}/(P_f + P_d)$ , because the antenna loss is included.

As for other method of calculating  $Q$  factors, following two methods are useful. In the reference [5], the lowest  $Q$  value of the next expression is given by J.S. McLearn.

$$Q_{lowest} = \frac{1}{k^3 a^3} + \frac{1}{ka} \quad (15)$$

In the reference [6], the relation of  $Q$  factor and bandwidth is given by the next expression.

$$Q_A = \frac{S-1}{\sqrt{S}} \cdot \frac{f_0}{f_B} \quad (16)$$

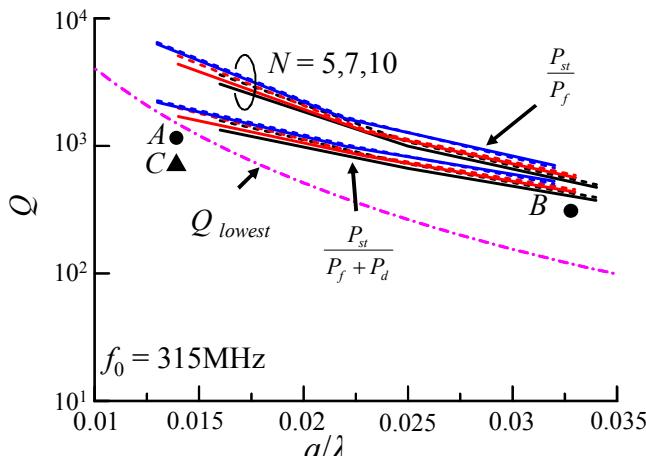


Fig. 9 Comparisons of  $Q$  factors

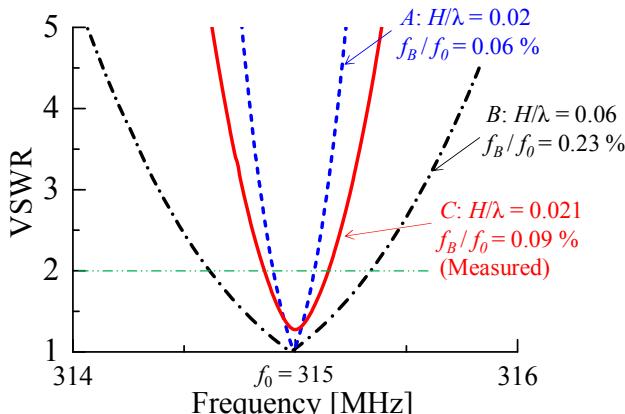


Fig. 10 Bandwidth characteristics

Here,  $S$  indicates the VSWR value.  $f_0$  and  $f_B$  indicate the center frequency and bandwidth, respectively.

Comparisons of  $Q$  factors based on Eq. (14) to Eq. (16) are shown in Fig. 9. The  $a$  at the horizontal axis indicates the radius of a sphere shown in Fig. 1. So, the twice value of  $a$  corresponds to  $H_A$ . The broken line indicates the result of Eq. (15). Black circles indicate the results of Eq. (16). The black triangle indicates a measured result. In calculation of black circles and triangle, the bandwidths shown in Fig. 10 are used.

It is interesting that black circles and triangle agree rather well with  $P_{st}/(P_f + P_d)$  results. Moreover, at small antenna structures near  $a/\lambda = 0.015$ , black circle,  $P_{st}/(P_f + P_d)$  and Eq. (15) results agree well.

As a result, adequateness of  $P_{st}$  results shown in Fig. 6 is ensured.

#### V. CONCLUSION

As a study object, normal-mode helical antennas are employed. From calculated results of electric and magnetic fields around antennas, stored electric and magnetic powers are obtained. The equivalence of two components at the self-resonant structures are shown numerically. The values of stored powers are shown to become 500 W to 2000 W at small antenna sizes. From stored power, antenna  $Q$  factors are estimated. As a result of good agreements of  $Q$  factors of three methods, adequateness of stored power results are ensured.

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